

A Perturbative Study of the Chirally Rotated Schrödinger Functional in QCD



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Stefan Sint

Overview:

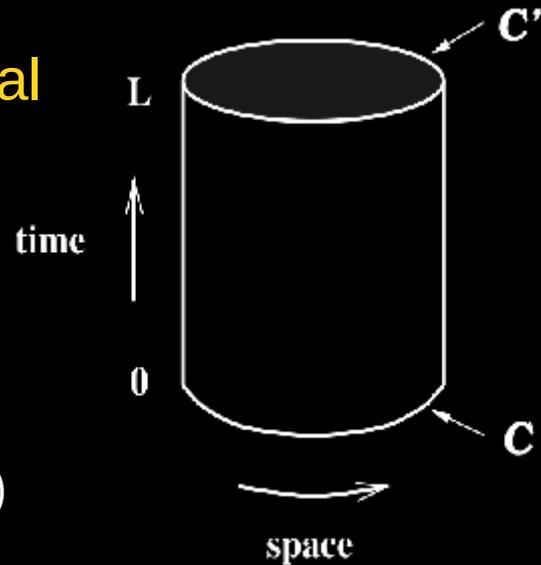
- Introduction: Schrödinger Functional schemes and automatic $O(a)$ improvement.
- The Chirally Rotated Schrödinger Functional (χ SF)
 - Definition
 - Renormalization & Improvement
 - Correlation Functions
- A perturbative study
 - Determination of coefficients
 - Checks of automatic $O(a)$ improvement.
- Applications
- Concluding Remarks



Introduction:

- Finite Volume (FV) schemes based on the Schrödinger Functional widely used in non perturbative renormalization.
- Hypercylindrical Euclidean manifold with temporal boundaries. [Lüscher et al. '92]

$$\mathcal{Z}[C, C'] = \int \mathcal{D}[A, \psi, \bar{\psi}] e^{-S[A, \psi, \bar{\psi}]}$$



- Boundary conditions:

-Gauge fields $C_k = \frac{i}{L} \text{diag}(\phi_{1k}, \dots, \phi_{Nk})$

$$P_+ \psi |_{x_0=0} = P_- \psi |_{x_0=T} = 0$$

-Fermion fields $\bar{\psi} P_- |_{x_0=0} = \bar{\psi} P_+ |_{x_0=T} = 0$

- Successfully applied in several renormalization problems:
(coupling in QCD, running quark masses, BSM, composite operators, ...)

Introduction:

- In any SF formulation, there are extra sources of cutoff effects generated at the boundaries. (Extra dim 4 operators localized at the boundaries).
- These can be removed through Symanzik improvement.
- For Wilson Fermions and gauge action, $O(a)$ effects are due to dim 5 operators in the bulk and dim 4 at the boundaries.
- $O(a)$ improvement is achieved by adding

-Bulk: (dim 5) $\bar{\psi} i \sigma_{\mu\nu} F_{\mu\nu} \psi \longrightarrow c_{\text{SW}}$

-Boundaries: (dim 4)

-Gauge: $\text{tr} \{ F_{kl} F_{kl} \} \longrightarrow c_t$

-Fermion: $\bar{\psi} P_{\pm} D_0 \psi \longrightarrow \tilde{c}_t$

Automatic $O(a)$ improvement

If \mathcal{R}_5 is a symmetry of the massless continuum theory,

$$\mathcal{O}_{\text{even}} \longrightarrow O(1), O(a^2), \dots$$

$$\mathcal{O}_{\text{odd}} \longrightarrow O(a), O(a^3), \dots$$

$$\mathcal{R}_5 : \psi \rightarrow i\gamma_5\psi$$
$$\bar{\psi} \rightarrow \bar{\psi}i\gamma_5$$



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But... $P_{\pm}\gamma_5 = \gamma_5 P_{\mp}$

Could we $[\tilde{\gamma}_5, \tilde{P}_{\pm}] = 0$

Yes !!!!

$$\tilde{\gamma}_5 = \gamma_5\tau_1 \quad \tilde{Q}_{\pm} = \frac{1}{2}(1 \pm i\gamma_0\gamma_5\tau_3)$$

Chirally Rotated SF

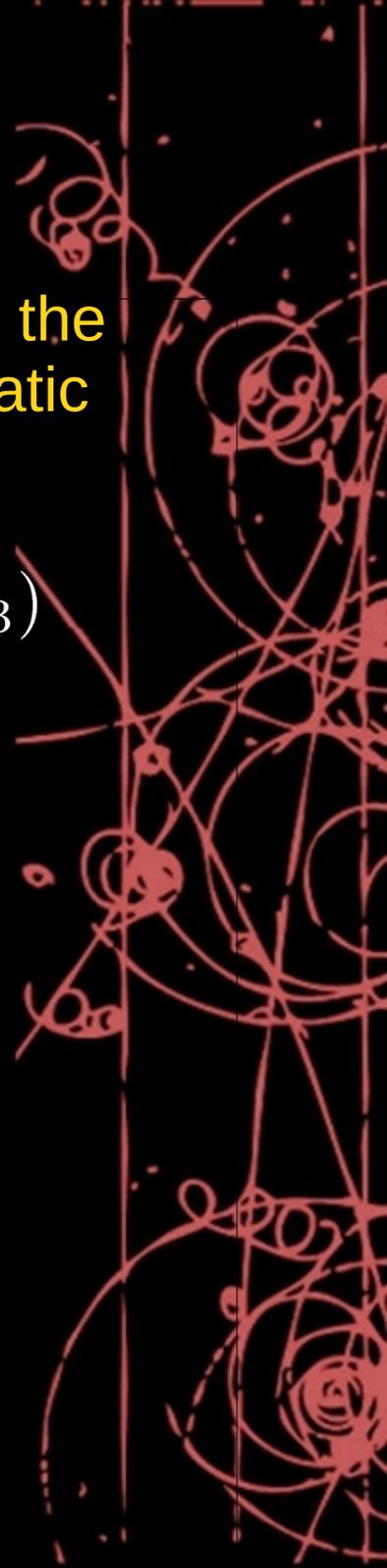
χ SF [Sint '05-'10]

$$\tilde{Q}_+ \psi \Big|_{x_0=0} = \tilde{Q}_- \psi \Big|_{x_0=T} = 0$$

$$\bar{\psi} \tilde{Q}_+ \Big|_{x_0=0} = \bar{\psi} \tilde{Q}_- \Big|_{x_0=T} = 0$$

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SF and χ SF related through a chiral rotation (Identical in the continuum and chiral limits).

$$\psi \longrightarrow R(\alpha)\psi$$

$$R(\alpha) = e^{i\gamma_5\tau^3\alpha/2} \quad \alpha = \pi/2$$

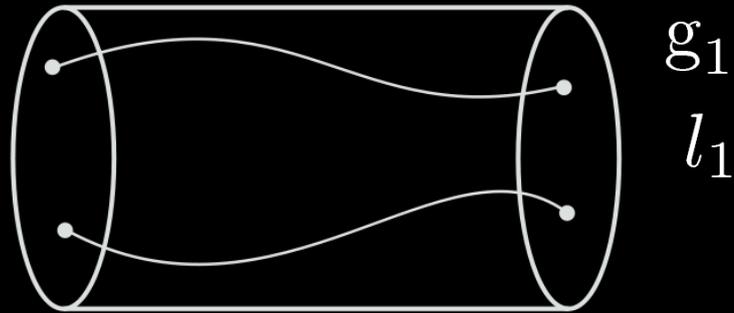
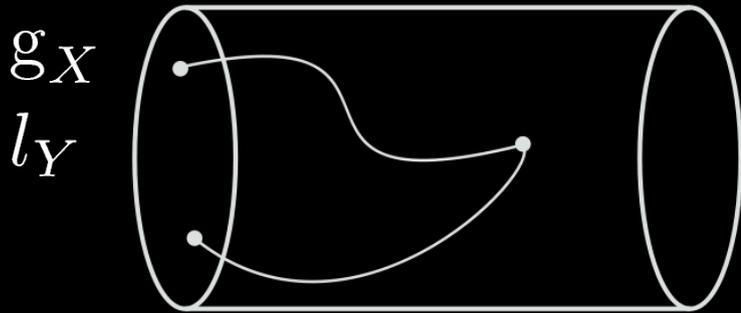
$$\bar{\psi} \longrightarrow \bar{\psi}R(\alpha)$$

Automatic $O(a)$ improvement is expected after renormalization and $O(a)$ imp (of boundaries):



χ SF correlation functions

- Boundary to bulk and boundary to boundary



$$g_X^{f_1 f_2}(x_0) = -\frac{1}{2} \langle X^{f_1 f_2}(x_0) Q_5^{f_2 f_1} \rangle$$

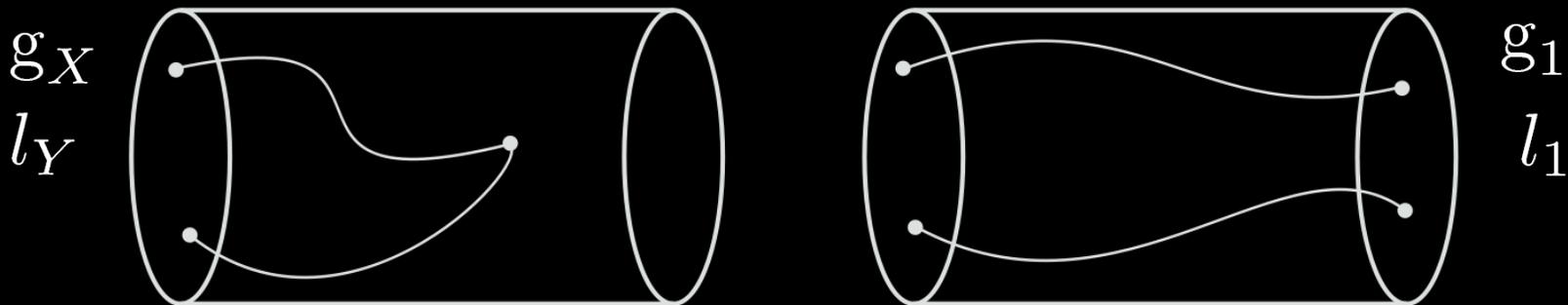
$$X = A_0, V_0, S, P$$

$$g_1^{f_1 f_2} = -\frac{1}{2} \langle Q_5^{f_1 f_2} Q_5^{\prime f_2 f_1} \rangle$$



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- These are related to standard SF correlation functions through the chiral twist

$$\langle O[R(\pi/2)\psi, \bar{\psi}R(\pi/2)] Q_5^{f_1 f_2} \rangle_{\chi SF} = \langle O[\psi, \bar{\psi}] O_5^{f_1 f_2} \rangle_{SF}$$

$$R(\alpha) = e^{i\gamma_5 \tau^3 \alpha/2}$$

χ SF correlation functions

Dictionary between SF and cSF correlation functions

$$\left. \begin{aligned} f_A &= g_A^{uu} = g_A^{dd} = -ig_V^{ud} = ig_V^{du} \\ f_P &= ig_S^{uu} = -ig_S^{dd} = g_P^{ud} = g_P^{du} \end{aligned} \right| \text{even}$$

$$\text{odd} \left| \begin{aligned} f_V &= g_V^{uu} = g_V^{dd} = -ig_A^{ud} = ig_A^{du} \\ f_S &= ig_P^{uu} = -ig_P^{dd} = g_S^{ud} = g_S^{du} \end{aligned} \right.$$

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Renormalization and improvement

Renormalization and improvement conditions evaluated order by order in perturbation theory.

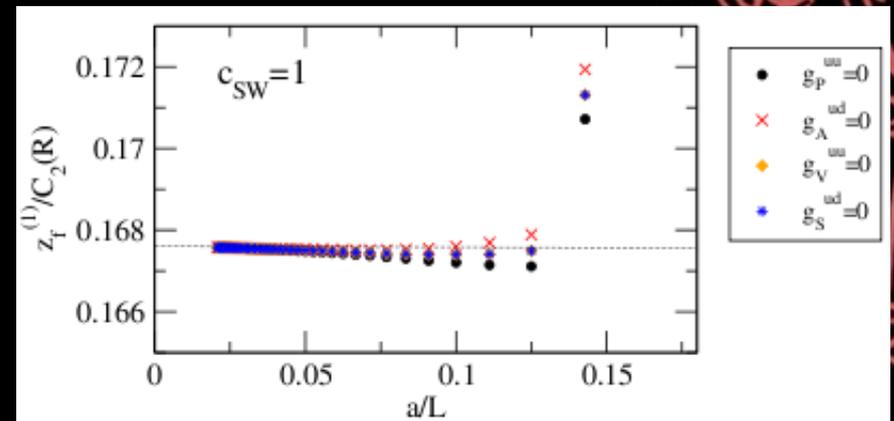
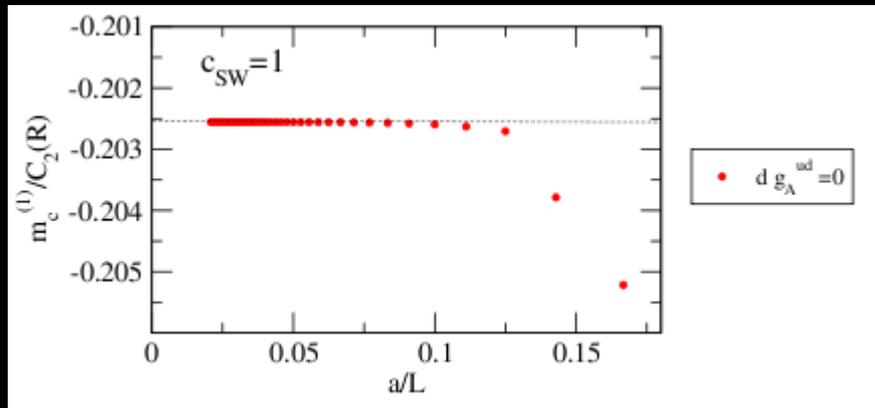
- m_c : $m_{\text{PCAC}} = 0$
- z_f : require a \mathcal{P}_5 odd observable to vanish.
 $g_A^{ud} = 0$; $g_V^{uu} = 0$; $g_P^{uu} = 0$; $g_S^{ud} = 0$
- d_s : require absence of $O(a)$ effects in \mathcal{P}_5 even observable
- c_t : require absence of $O(a)$ terms in the 1-loop coupling

Renormalization and improvement

For m_c and z_f :

$$m_c^{(1)} : \begin{cases} -0.2025565(1) \times C_2(\mathcal{R}), & c_{\text{SW}} = 1 \\ -0.325721(7) \times C_2(\mathcal{R}), & c_{\text{SW}} = 0, \end{cases}$$

$$z_f^{(1)} : \begin{cases} 0.167572(2) \times C_2(\mathcal{R}), & c_{\text{SW}} = 1 \\ 0.33023(6) \times C_2(\mathcal{R}), & c_{\text{SW}} = 0, \end{cases}$$



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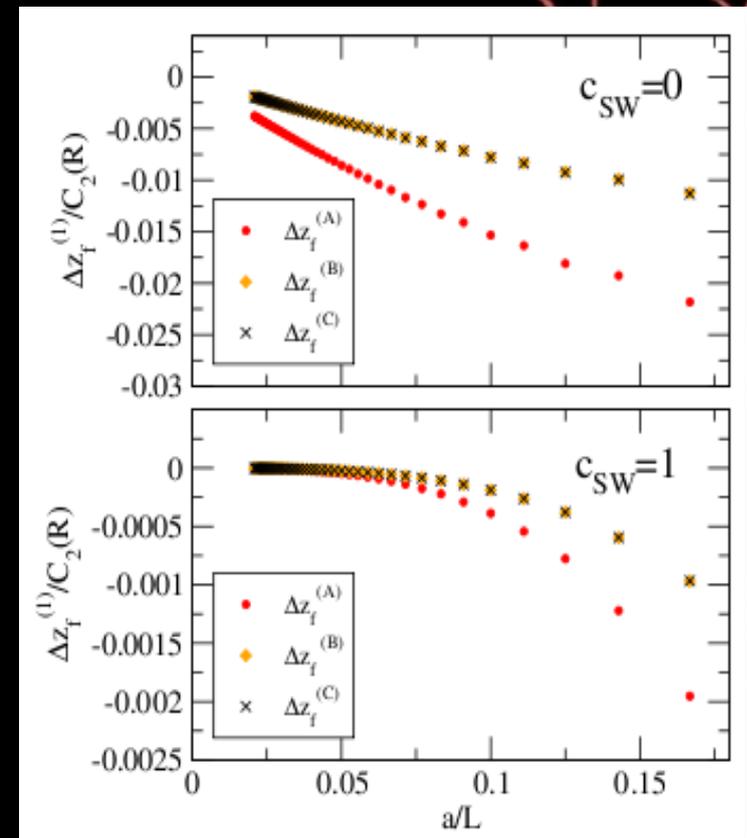
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Different renormalization conditions

→ $O(a)$ differences in z_f

$$\begin{cases} \Delta z_f^{(A)} = z_f^{(1)} \left|_{g_A^{ud}} - z_f^{(1)} \right|_{g_P^{uu'}} \\ \Delta z_f^{(B)} = z_f^{(1)} \left|_{g_V^{uu'}} - z_f^{(1)} \right|_{g_P^{uu'}} \\ \Delta z_f^{(C)} = z_f^{(1)} \left|_{g_S^{ud}} - z_f^{(1)} \right|_{g_P^{uu'}} \end{cases}$$



Renormalization and improvement

For d_s we demand $O(a)$ effects to be absent from the ratio
(several θ)

$$\frac{[g_P^{ud}(x_0, \theta, a/L)]_R}{[g_P^{ud}(x_0, 0, a/L)]_R} \Big|_{x_0=T/2} \longrightarrow d_s^{(1)} = -0.0009(3) \times C_2(\mathcal{R})$$

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For c_t we demand $O(a)$ effects to be absent from the SF coupling (standard definition through a background field).

$$\frac{1}{\bar{g}^2} = \frac{\partial \Gamma / \partial \eta |_{\eta=0}}{\partial \Gamma_0 / \partial \eta |_{\eta=0}}$$

$$c_t^{(1,1)} : \begin{cases} 0.006888(3), & c_{\text{SW}} = 1 \\ -0.00661445(5), & c_{\text{SW}} = 0, \end{cases}$$

Check automatic $O(a)$ improvement

i) Correct realization of the boundary conditions

Boundary bilinears defined with opposite projectors

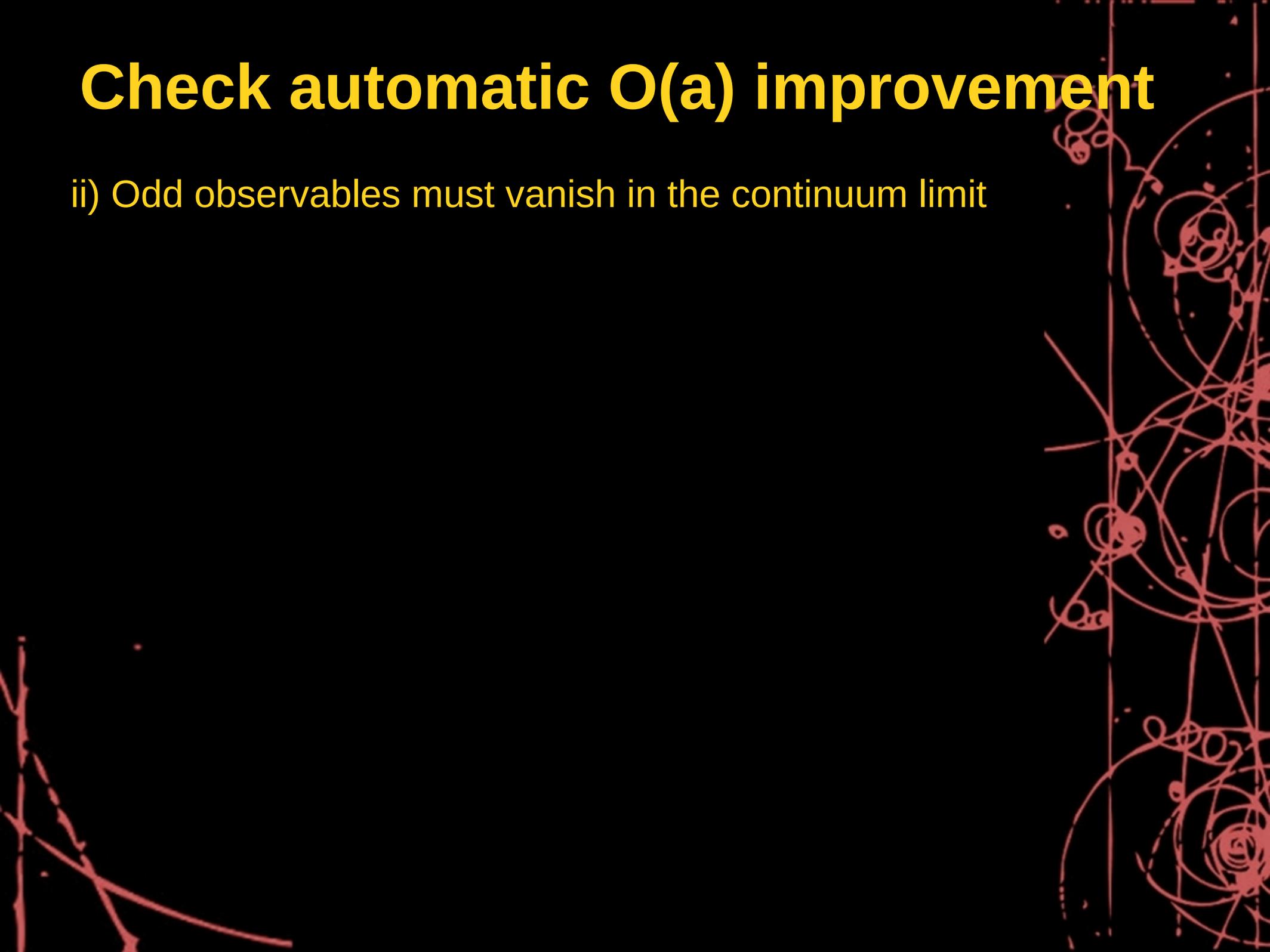
$$Q_5 \sim \bar{\zeta} \Gamma \tilde{Q}_{\pm} \zeta$$

$$\tilde{Q}_{\mp}$$

$$S(x, y) \bar{Q}_+ |_{y_0=0} = 0$$

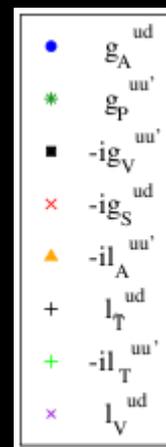
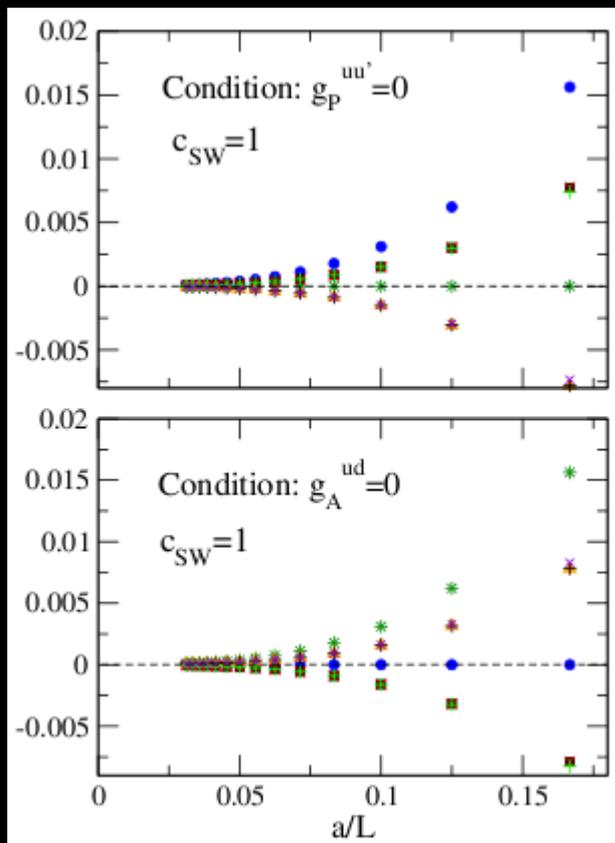
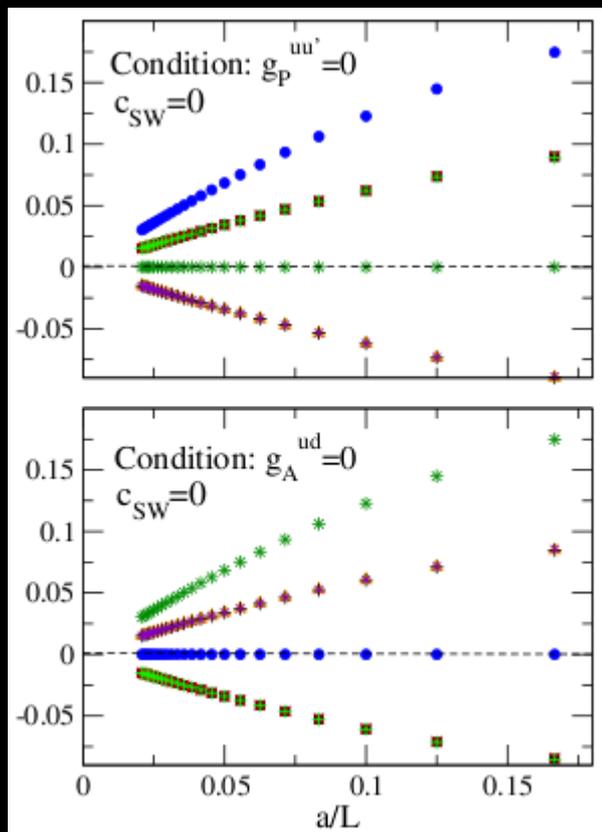
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ii) Odd observables must vanish in the continuum limit



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Check automatic $O(a)$ improvement

iii) Universality between SF and χ SF: Ratios between renormalized correlation functions converge.

$$\frac{g_A^{uu'}(T/2)/\sqrt{g_1}}{f_A(T/2)/\sqrt{f_1}} \longrightarrow 1 \qquad \frac{g_P^{ud}(T/2)/\sqrt{g_1}}{f_P(T/2)/\sqrt{f_1}} \longrightarrow 1$$

Check automatic $O(a)$ improvement

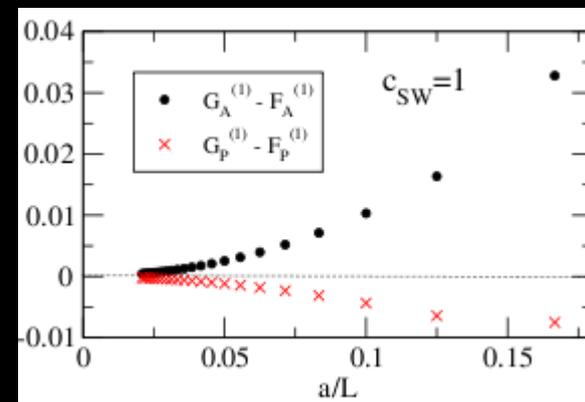
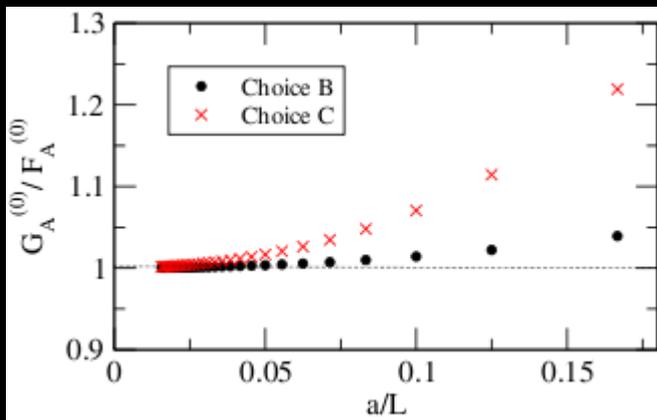
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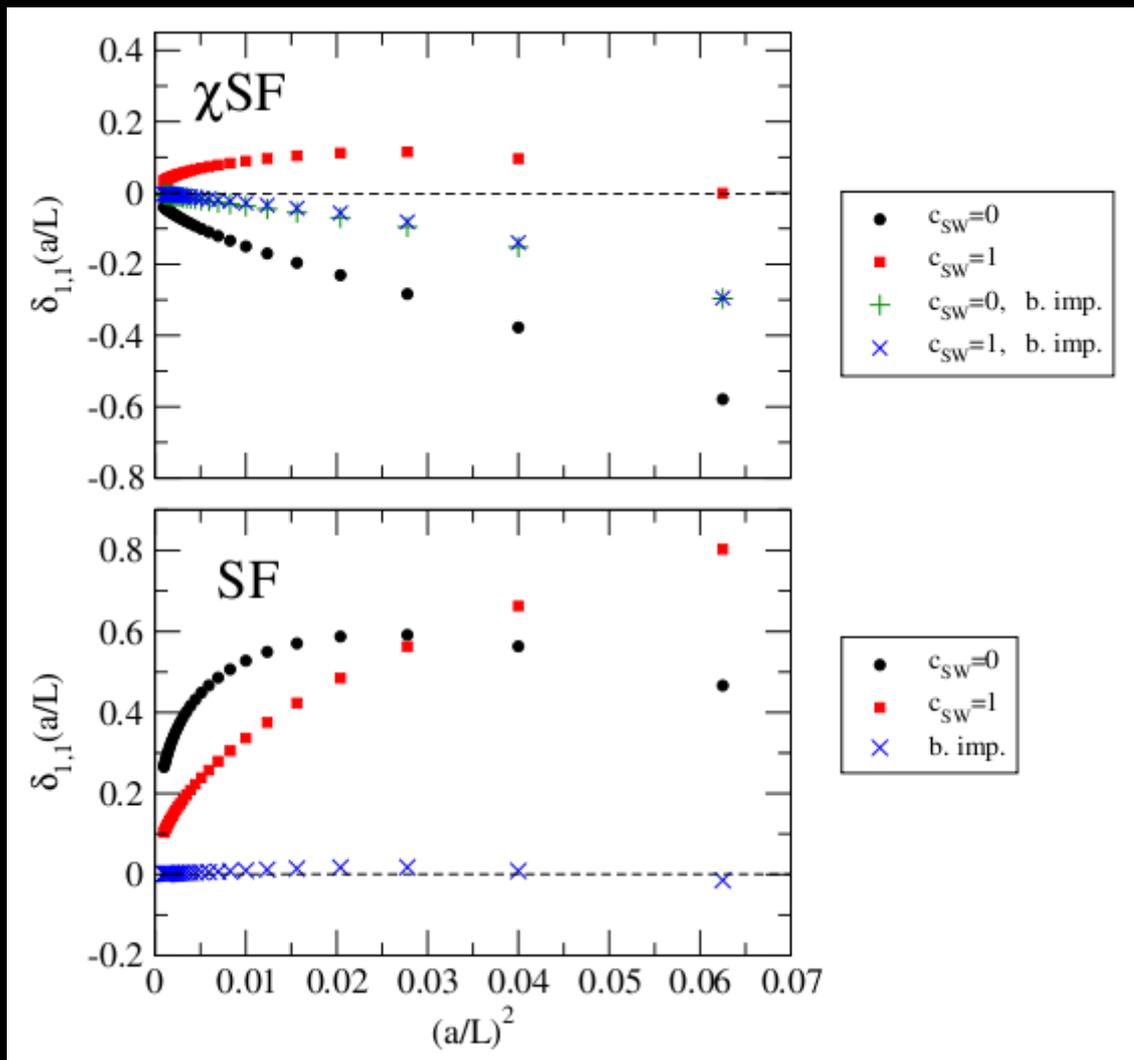
Tree-level;

1-loop:



Setup ready!!!!

Cutoff effects in the SSF



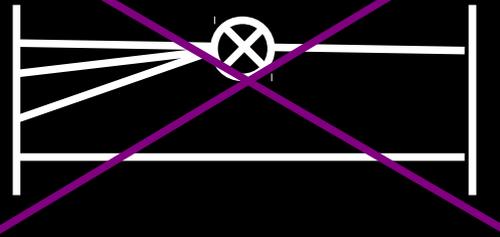
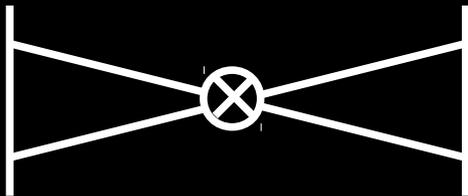
[S. Sint, P.V. '11-'12]

Applications:

- Computation of finite renormalization factors (see Mattia Dalla Brida's talk) [B. Leder, S. Sint, '10]

$$Z_A = i \frac{g_{\tilde{V}}^{ud}(T/2)}{g_A^{uu}(T/2)} \quad Z_V = \frac{g_{\tilde{V}}^{ud}(T/2)}{g_V^{ud}(T/2)}$$

- Renormalization of 4-fermion operators:
Simpler observables & no operator improvement



- Twist-2 operators. [J. Gonzalez Lopez et al, '12]

Conclusions

- The χ SF implements the mechanism of automatic $O(a)$ improvement in the SF setup.
- We have determined to 1-loop in PT the necessary coefficients for the renormalization and $O(a)$ improvement of the setup: $m_c^{(1)}$, $z_f^{(1)}$, $d_s^{(1)}$, $c_t^{(1)}$
- We have confirmed that after fixing these parameters, automatic $O(a)$ improvement is at hold (at 1-loop).
- The running coupling has been computed to 1-loop in pt.
- See Mattia's talk for a $N_f=2$ dynamical calculation.

..... thank you very much !!!

